Prof. Hans Peter Büchler SS 2011, 30th of November 2011

## 1. Stark effect in harmonic oscillator (Oral)

The Hamiltonian of a 1-D harmonic oscillator in a homogeneous electric field E is

$$H = \frac{1}{2}(P^2 + Q^2) + eEQ.$$
 (1)

We consider the second part of the hamiltonian as a pertubation of the free oscillator  $H_1 := Q$ , with  $\lambda := eE$ .

- (a) Calculate the pertubed eigenvalues and eigenfunction up to the second order in  $\lambda$ .
- (b) Compare the perturbative result with the exact solution.

## 2. Spin orbit and Zeeman effect (Homework)

We consider a system with angular momentum L = 1 and spin S = 1/2. Within  $H_0$  all these spin states are degenerate. Now, we add the spin-orbit coupling  $H_{so} = \gamma \mathbf{SL}$  as a perturbation.

- (a) Use the degenerate perturbation theory to determine the splitting of the degeneracy under the spin orbit coupling. (Hint : use the addition of angular momentum to find a appropriate basis. Note that  $\mathbf{SL} = (\mathbf{J}^2 \mathbf{L}^2 \mathbf{S}^2)/2$  with  $\mathbf{J}$  the total angular momentum).
- (b) We now apply a magnetic field

$$H = H_0 + H_{so} + H_m, \tag{2}$$

where  $H_m = -\vec{\mu} \cdot \vec{B}$ . Using perturbation theory solve this hamiltonian for two different cases :

- (i) weak magnetic field (Zeeman effect), consider here a pertubation of the spin orbit levels.
- (ii) strong magnetic field (Paschen-Back effect), where the spin orbit coupling can be ignored.

## 3. Van-der-Waals interaction (Homework)

Consider two hydrogen atoms (A, B) separated by a large distance  $\mathbf{R}$ , so that there is no overlap of electronic charge distribution, and have fixed nuclei (Born-Oppenheimer approximation). Let  $H_{\alpha}$  be the hamiltonian of a single hydrogen atom. We can then write the Hamiltonian of the full system as

$$H = H_A + H_B + H_{int},\tag{3}$$

where  $H_{int}$  describes the Coulomb interaction between the electrons and the positively charged nuclei, i.e.,

$$H_{int} = \frac{e^2}{R} + \frac{e^2}{|\mathbf{R} + \mathbf{r}_B - \mathbf{r}_A|} - \frac{e^2}{|\mathbf{R} + \mathbf{r}_B|} - \frac{e^2}{|\mathbf{R} - \mathbf{r}_A|},$$
(4)

with  $\mathbf{r}_i$   $(i \in A, B)$  the position of an electron with respect to the nucleus *i*. For large  $\mathbf{R} \gg a_0$  (Bohr-radius) one can develop  $H_{int}$  in powers of  $\mathbf{r}_i/\mathbf{R}$  and obtain the dipole-dipole interaction

$$H_{int} = \frac{1}{R^3} \left[ \boldsymbol{d}_A \boldsymbol{d}_B - \frac{3(\boldsymbol{d}_A \cdot \boldsymbol{R})(\boldsymbol{d}_B \cdot \boldsymbol{R})}{R^2} \right],$$
(5)

with the dipole operator  $d_i = er_i$ . In the following we want to obtain the energy correction of the states  $|g\rangle_A |g\rangle_B$  with the two atoms in the lowest energy state  $|g\rangle_i$ . We consider the distance **R** between the atoms as fixed.

- a) Show that the first order in perturbation theory gives no contribution.
- b) For the second order perturbation theory, we simplify the model, where we assume that the dipole operator couples only to excited state  $|e\rangle_i$ , i.e.,  $\langle e|\mathbf{d}_i|g\rangle_i = d/\sqrt{3}\mathbf{e}_z$  with  $\mathbf{e}_z$  the unit vector along the z-axis. Show that the energy correction is always negative in the second-order and thus describes an attraction between the two atoms. Comment on the spatial dependence of the resulting interaction. Show that this interaction goes like  $1/R^6$ .
- c) We now consider the degenerate excited p-states of the hydrogen atoms, i.e. , each atom has a 3-fold degenerate state  $|2, 1, m\rangle$  with principal quantum number n = 2, angular momentum l = 1, and  $m = 0, \pm 1$ . Show the anistropic behavior of the Van der Waals interaction.