

## Solid State Theory, Exercises VIII

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Prof. Hans Peter Büchler SS 2012, 29th of June 2012

### The RPA dielectric function

The general form of the screened susceptibility is given by

$$\chi_0(\mathbf{q}, \omega) = 2 \sum_{\mathbf{k}} \frac{n_{\mathbf{k}} - n_{\mathbf{k}+\mathbf{q}}}{\omega - \epsilon_{\mathbf{k}+\mathbf{q}} + \epsilon_{\mathbf{k}} + i\eta} \quad (1)$$

(a) Show that the screened susceptibility is :

$$\begin{aligned} \text{Re}[\chi_0(\mathbf{q}, \omega)] = -g \left\{ \frac{1}{2} + \frac{k_F}{4q} \left[ \frac{(\omega + q^2/2m)^2}{(qv_F)^2} - 1 \right] \ln \left[ \frac{\omega - qv_F + q^2/2m}{\omega + qv_F + q^2/2m} \right] \right. \\ \left. - \frac{k_F}{4q} \left[ \frac{(\omega - q^2/2m)^2}{(qv_F)^2} - 1 \right] \ln \left[ \frac{\omega - qv_F - q^2/2m}{\omega + qv_F - q^2/2m} \right] \right\}, \end{aligned}$$

and

$$\text{Im}[\chi_0(\mathbf{q}, \omega)] = \begin{cases} -(1+g) \frac{\pi}{2} \frac{\omega}{qv_F} & \text{for } 0 < \omega < v_F q - q^2/2m \\ (1+g) \frac{\pi}{2} \frac{p_F}{q} \left[ 1 - \frac{(\omega - q^2/2m)^2}{(qv_F)^2} \right] & \text{for } v_F q - q^2/2m < \omega < v_F q + q^2/2m \\ 0 & \text{for } v_F q + q^2/2m < \omega < \infty, \end{cases} \quad (2)$$

where  $g$  is the density of states.

(b) Using the above formula compute the dielectric function  $\epsilon_{RPA}(\mathbf{q}, \omega)$ .

(c) Write down  $\text{Re}[\epsilon_{RPA}]$  in the static limit  $\omega \rightarrow 0$ , and sketch it as a function of  $q$ . Expand at small  $q$  to recover the Thomas-Fermi screening form. What happens at  $q = 2k_F$ ? Show that at this point the dielectric constant is continuous but non-differentiable. Explain why this a special value of  $q$ .

(d) Expand  $\text{Re}[\epsilon_{RPA}(\mathbf{q}, \omega)]$  at small  $q$  to order  $q^2$  (this requires expansion of the logarithms to quite a high order in  $q/\omega$ ). You should recover the plasmon dispersion given in exercise sheet VII.

Solutions due on : 6th of July, 2012