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Exercise 1: Current conservation (2 points, Oral)

Consider the continuity equation in real space,

$$\partial_t \rho(x) = -\frac{i}{\hbar} [\rho(x), H], \quad (1)$$

and a general Hamiltonian of interacting fermions in second quantized form,

$$H = \int d^3x \frac{\hbar^2}{2m} \nabla \psi^\dagger(x) \nabla \psi(x) + \frac{1}{2} \int d^3x d^3y \psi^\dagger(x) \psi^\dagger(y) \psi(y) \psi(x) V(x-y), \quad (2)$$

to show that

$$\partial_t \rho(x) = -\nabla j(x), \quad (3)$$

where the current is

$$j(x) = -\frac{i}{2m} [(\nabla \psi^\dagger(x)) \psi(x) - \psi^\dagger(x) \nabla \psi(x)], \quad (4)$$

and the density operator in terms of fields was defined in the lecture, $\rho(x) = \psi^\dagger(x) \psi(x)$.

Exercise 2: Sum rules (5 points, Oral)

In this exercise, we are going to prove the sum rule of the structure factor, $S(\mathbf{q}, w)$,

$$\int_0^\infty dw w S(q, w) = \frac{N}{V} \frac{q^2}{2m}, \quad (5)$$

implying the following limit of the density-correlator,

$$\chi(q, w \rightarrow \infty) = \frac{N}{V} \frac{q^2}{2m}. \quad (6)$$

a) Fourier transform $j(x)$ to obtain the current in the momentum space,

$$j_q = \sum_{k, \sigma} [(k + q/2)/m] c_{k, \sigma}^\dagger c_{k+q, \sigma}. \quad (7)$$

b) Show that, in momentum space, the continuity equation (1) leads to,

$$[\rho_q, H] = -i\hbar q \cdot j_q. \quad (8)$$

c) Show that

$$[[\rho_q, H], \rho_{-q}] = [q \cdot j_q, \rho_{-q}] = \frac{Nq^2}{m}. \quad (9)$$

d) Evaluate Eq. (9) in the ground-state, i.e. the Fermi sea $|\Psi_0\rangle$, to show,

$$\langle \Psi_0 | [[\rho_q, H], \rho_{-q}] | \Psi_0 \rangle = \frac{2}{V} \int_0^\infty dw w S(q, w). \quad (10)$$

For this purpose, first use the first commutator equality of task b) and introduce the identity, $\sum_n |n\rangle \langle n|$,

$$\langle \Psi_0 | [[\rho_q, H], \rho_{-q}] | \Psi_0 \rangle = \sum_n \langle 0 | q \cdot j_q | n \rangle \langle n | \rho_{-q} | 0 \rangle - \sum_n \langle 0 | \rho_{-q} | n \rangle \langle n | q \cdot j_q | 0 \rangle. \quad (11)$$

Then show that $\langle n | q \cdot j_q | 0 \rangle = -w_{n0} \langle n | \rho_q | 0 \rangle$, and convert the sum into a continuous integral in w .

e) Use this result to show that

$$\chi(q, w \rightarrow \infty) = \frac{2}{w^2} \int_0^\infty dw' w' S(q, w') = \frac{N}{V} \frac{q^2}{mw^2}. \quad (12)$$

Exercise 3: Response function of the free Fermi gas (2 points, Written)

Consider the response function seen in the course,

$$\chi(q, w) = \frac{1}{V\hbar} \sum_n |\langle n | \rho_{-q} | \Psi_0 \rangle|^2 \left(\frac{1}{w - w_{n0} + i\eta} - \frac{1}{w + w_{n0} + i\eta} \right). \quad (13)$$

Show that, for the free Fermi gas, it can be written as

$$\chi(q, w) = 2 \int \frac{dk}{(2\pi)^3} \frac{n_k - n_{k+q}}{\hbar w - (\varepsilon_{k+q} - \varepsilon_k) + i\eta}, \quad (14)$$

where the energies are $\varepsilon_k = \hbar^2 k^2 / 2m$, and n_k is the Fermi-Dirac distribution.

Hint: Consider the possible excitation states $|n\rangle$ of a free Fermi gas, i.e. particle-hole excitations.