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You can find detailed information about the lecture and the exercises on the website <http://www.itp3.uni-stuttgart.de/lehre/vorlesungen/QM-main.html>. There are two types of exercises: *Written* exercises have to be handed in and will be corrected by the tutors. *Oral* exercises have to be prepared for the exercise session and will be presented by one of the students. To be admitted to the exam, we require (a) 80% of the written exercises to be solved or treated adequately, (b) 66% of the oral exercises to be prepared and (c) two exercises to be presented on the black board.

### Exercise 1: Plane waves (Written, 4 points)

In an interval  $[0, L]$  with periodic boundary conditions, plane waves are given by

$$\psi_n = \frac{1}{\sqrt{L}} \exp\left(\frac{i}{\hbar} p_n x\right), \quad \text{where } p_n = \frac{2\pi\hbar n}{L}. \quad (1)$$

- a) Show that they form a complete set of orthonormal basis functions. In other words, show that:

$$\int_0^L dx \psi_n^*(x) \psi_m(x) = \delta_{n,m} \quad (\text{orthonormality}), \quad (2)$$

$$\sum_{n=-\infty}^{\infty} \psi_n^*(x) \psi_n(x') = \delta(x - x') \quad (\text{completeness}). \quad (3)$$

- b) In the limit  $L \rightarrow \infty$ , the orthogonality and completeness relations read

$$\int_{-\infty}^{\infty} dx \exp\left[-\frac{i}{\hbar}(p - p')x\right] = 2\pi\hbar \delta(p - p') \quad (\text{orthogonality}), \quad (4)$$

$$\int_{-\infty}^{\infty} \frac{dp}{2\pi\hbar} \exp\left[\frac{i}{\hbar}(x - x')p\right] = \delta(x - x') \quad (\text{completeness}). \quad (5)$$

Perform the limit  $L \rightarrow \infty$  explicitly for  $\psi_n$  by using an appropriate prefactor and show both the orthogonality (4) and the completeness (5). (Note that the basis functions are no longer normalizable and are thus called a *general set of basis functions*.)

**Exercise 2: Classical action (Written, 5 points)**

In classical mechanics, the action  $S$  is defined as

$$S[q](t) = \int_{t_0}^t dt' L(\dot{q}(t'), q(t'), t'), \quad (6)$$

where  $L$  is the *Lagrangian function* (also called *Lagrangian*) and  $q$  is some generalized coordinate of the system.

a) Calculate the action  $S(q_\tau, \tau)$  along classical paths for the following systems:

- free particle,
- harmonic oscillator with  $V(q) = \frac{1}{2}m\omega^2 q^2$ ,
- constant force  $F$ ,

where  $q(t = \tau_0) = 0$  is the starting point and  $q(t = \tau) = q_\tau$  is the endpoint of the path at some time  $\tau > \tau_0$ .

b) For classical paths with a fixed starting point  $q(\tau_0) = q_0$ , one can interpret  $S$  as a function of  $q_\tau$  and  $\tau$ , i.e.  $S = S(q_\tau, \tau)$ . Show that

$$\frac{\partial S}{\partial q_\tau} = p_\tau, \quad \frac{\partial S}{\partial \tau} = -H, \quad (7)$$

where  $H$  is the Hamiltonian function of the system.

*Hint:* For some fixed starting point  $q_0$ , each path can be parametrized by some arbitrary point  $q_\tau$  on that path such that  $q(t)$  itself can be thought of as a function of both  $q_\tau$  and some general time  $t$ .

**Exercise 3: Free propagator (Oral)**

The propagator  $K(x, x', t)$  of a Hamiltonian operator  $H$  (some differential operator in  $x$ ) is defined as the solution of the Schrödinger equation

$$i\hbar\partial_t K(x, x', t) = HK(x, x', t) \quad (8)$$

with the initial condition  $K(x, x', 0) = \delta(x - x')$ .

a) Show that for the initial condition  $\psi(x, t = 0) = \psi_0(x)$  the solution of the Schrödinger equation is given by

$$\psi(x, t) = \int dx' K(x, x', t)\psi_0(x'). \quad (9)$$

b) Show using Fourier transformation (plane wave expansion) that the propagator of a free particle with Hamiltonian operator  $H = -\hbar^2\partial_x^2/2m$  is given by

$$K(x, x', t) = \sqrt{\frac{m}{2\pi\hbar it}} \exp\left(\frac{im(x - x')^2}{2\hbar t}\right). \quad (10)$$