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Bose-Einstein condensation of cold gases in traps

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Hauptseminar: Physik der kalten Gase

30.04.2013

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Figure : Density distribution of Sodium atoms with $T_c=2~\mu {\rm K}.$ From: Nobel lecture of Wolfgang Ketterle (2001)

					2 2 2	200 200 201)2)4 .1		Cs Cr D <u>y</u>	5								
Group — ↓ Period	• 1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18
1	H																	2 He
2	3 Li	4 Be											5 B	6 C	7 N	8 0	9 F	10 Ne
3	11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl	18 Ar
4	19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
5	37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 1	54 Xe
6	55 Cs	56 Ba		72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 TI	82 Pb	83 Bi	84 Po	85 At	86 Rn
7	87 Fr	88 Ra		104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt	110 Ds	111 Rg	112 Cn	113 Uut	114 Fl	115 Uup	116 Lv	117 Uus	118 Uuo
				_	_								_					
	La	nthan	ides	57 La	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu
		Actin	ides	89 Ac	90 Th	91 Pa	92 U	93 No	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Ur

Rb Na Li

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1995 1998

2001

Figure : Time line and periodic table of BEC.

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Outline				

Ideal bose gas

- Harmonic potential
- Critical temperature T_c
- Density distribution

2 Experimental techniques

- Typical setup
- Radiation forces on atoms
- Laser cooling
- Evaporative cooling

3 Weakly interacting bose gas

- Thomas-Fermi approximation
- Healing length

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Harmon	ic potential			

Isotropic harmonic potential

$$V_{\rm ho}(\mathbf{r}) = \frac{1}{2}m\omega_{\rm ho}^2\mathbf{r}^2$$

Oscillator length

$$a_{\rm ho} = \sqrt{\frac{\hbar}{m\omega_{\rm ho}}} \approx 5\,\mu{\rm m}$$



Figure : Potential and ground state of a harmonic oscillator.

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Ideal bo	se gas			

Hamiltonian

$$H = \sum_{i} H_{i} = \sum_{i} \left[-\frac{\hbar^{2}}{2m} \nabla_{i}^{2} + V_{\rm ho}(\mathbf{r}_{i}) \right]$$

Single-particle energy and ground state ($n_x = n_y = n_z = 0$)

$$E = \hbar\omega_{\rm ho} \left(n_x + n_y + n_z + \frac{3}{2} \right)$$

$$\Phi(\mathbf{r}) = \left(\frac{m\,\omega_{\rm ho}}{\pi\,\hbar}\right)^{3/4} \exp\left[-\frac{m}{2\hbar}\omega_{\rm ho}\,r^2\right] = \left(\frac{1}{\sqrt{\pi}\,a_{\rm ho}}\right)^{3/2} \exp\left[-\frac{r^2}{2a_{\rm ho}^2}\right]$$

N-particle wave function and density distribution

$$\Psi(\mathbf{r}_1, \dots, \mathbf{r}_N) = \prod_i \Phi(\mathbf{r}_i) \qquad \qquad n(\mathbf{r}) = N \left| \Phi(\mathbf{r}) \right|^2$$

Blackboard: Density of states

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Densitv	of states			

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Critical	temperature	$e T_c$		

Bose distribution and density of states

$$f(E,T) = \frac{1}{e^{(E-\mu)/kT} - 1} \qquad \qquad g(E) = \frac{1}{2(\hbar\omega_{\rm ho})^3} E^2$$

Particle number in excited states, zero-point energy neglected

$$N_{\rm ex} = \int_0^\infty \mathrm{d}E \ g(E) f(E,T)$$

All particles in excited states at T_c

$$N = N_{\text{ex}}(T = T_c, \mu = 0) = \int_0^\infty dE \, \frac{g(E)}{e^{E/kT_c} - 1}$$
$$= \frac{1}{2} \left(\frac{kT_c}{\hbar\omega_{\text{ho}}}\right)^3 \int_0^\infty dz \, \frac{z^2}{e^z - 1} = \frac{1}{2} \left(\frac{kT_c}{\hbar\omega_{\text{ho}}}\right)^3 \Gamma(3)\zeta(3)$$
$$= \zeta(3) \left(\frac{kT_c}{\hbar\omega_{\text{ho}}}\right)^3$$







Figure : Noninteracting bosons in a spherical trap at temperature $T = 0.9 T_c$, Length z in units of $a_{\rm ho}$. From: Dalfovo et. al. (1999)

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Conclusion

Density distribution



Figure : Density distribution of Sodium atoms with $T_c=2~\mu{\rm K}.$ From: Nobel lecture of Wolfgang Ketterle (2001)

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Typical setup



Figure : Setup of a BEC experiment. From: Nobel lecture of Wolfgang Ketterle (2001)





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Dinole	force			

Stark effect

$$V_S = -\frac{1}{2}\alpha'(\omega)\left\langle E^2\right\rangle_t$$

Dipole force

$$\begin{aligned} \mathbf{F}_{\text{dip}} &= -\nabla V_S(\mathbf{r}) \\ &= \frac{1}{2} \alpha'(\omega) \, \nabla \left\langle E^2 \right\rangle_t \end{aligned}$$



Figure : Atoms in an optical lattice potential. From: Wikipedia - Optisches Gitter



Figure : Laser beam profile and potential of an atom in this beam. From: Diploma thesis at PI5



Figure : Real (red) and imaginary part (green) of the polarizability over frequency.

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Scatteri	ng force			

Photon momentum

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$$\mathbf{p}=\hbar\mathbf{q}$$

Scattering force

$$\mathbf{F}_{s}(\omega-\omega_{0})=\mathbf{\dot{p}}=\hbar\mathbf{q}\,\Gamma_{g}(\omega-\omega_{0})$$

Absorption rate

$$\Gamma_{\rm g}(\omega-\omega_0) = -\frac{1}{2}\alpha''(\omega-\omega_0)\left\langle E^2\right\rangle_t$$

Recoil limit

$$kT_r = \frac{\hbar^2 q^2}{2m} \quad \Rightarrow \quad T_r \approx 1 \,\mu\mathrm{K}$$



Figure : Absorption of a photon followed by spontaneous emission in a random direction.



Figure : Real (red) and imaginary part (green) of the polarizability over frequency.

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Zeeman	slower			

Doppler effect

$$\omega = \omega_0 - kv$$

Magnetic field cancels decreasing Doppler shift

$$\begin{split} \hbar \omega &= \hbar \omega_0 - V_D + V_z \\ &= \hbar \omega_0 - \hbar k v + g m \mu_B B \end{split}$$

Constant deceleration of atoms on resonance

$$-\dot{v} = a = \frac{F_s}{m} \approx 10^5 \frac{m}{s^2}$$

$$v_0^2 - v^2 = 2az \Rightarrow v(z) = v_0 \sqrt{1 - \frac{2a}{v_0^2} z}$$



Figure : An atom moving towards a laser beam



Figure : Zeeman slower. From: Foot - Atomic Physics

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Doppler	cooling			

Doppler effect

$$\omega = \omega_0 - kv$$

 \Rightarrow Use red-detuned lasers ($\omega < \omega_0$)

Force of two counter-propagating laser beams

$$F = F_{s}(\omega - \omega_{0} + kv) -F_{s}(\omega - \omega_{0} - kv) \approx -\beta v$$

Atoms are cooled, but not trapped.

 \Rightarrow MOT: Add a spacial dependence



Figure : Moving atom in two counterpropagating laser beams.



Figure : Force over the velocity of an atom. From: Foot - Atomic Physics

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Magneto-optical trap

Magnetic quadrupole field: Linear Zeeman shift at origin

 $V_Z = gm\mu_B B \propto r$

Counter-propagating laser beams: Red-detuned, σ_- and σ_+ polarized



Figure : Trapped lithium atoms. From: Quantum optics group at ETH Zürich



Figure : Setup of a MOT. From: Foot - Atomic Physics



Magneto-optical trap



Figure : Schematic of a 1D MOT with a $J=0 \leftrightarrow J=1$ transition. From: Foot - Atomic Physics

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Magnet	ic tran			

Harmonic trap

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$$B = \frac{B''}{2}r^2 + B_0$$

Zeeman effect

$$V_Z(\mathbf{r}) = gm\mu_B B$$

Trapping force

$$\mathbf{F}(\mathbf{r}) = -\nabla V_Z \propto -gm \, \mathbf{r}$$

Low-field seekers (gm > 0) are trapped, high-field seekers (gm < 0) are repelled.



Figure : Currents in an loffe-Pritchard trap inducing a harmonic magnetic field. From: Pethick, Smith - BEC in dilute gases



Figure : Magnetic trap potential with a low-field seeking state (red) and a high-field seeking state (green).



Figure : Atoms in a harmonic potential. From: Pethick, Smith - BEC in dilute gases



Figure : (a) Maxwell-Boltzmann energy distribution at $T = T_0$, (b) Cut-off of "hot" atoms with $E > E_c$, (c) New equilibrium temperature $T_1 < T_0$. From: Diploma thesis at PI5.

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Thomas	-Fermi appr	oximation		

Gross-Pitaevskii equation

$$\left[-\frac{\hbar^2}{2m}\Delta + V_{\rm ho}(\mathbf{r}) + g|\Psi|^2\right]\Psi = \mu\Psi$$

Coupling constant $g=\frac{4\pi\hbar^2}{m}a>0$ for repulsive interactions.

Thomas-Fermi approximation

For $\frac{Na}{a_{\text{bo}}} \gg 1$ kinetic energy is small and can be neglected:

$$\left[V_{\rm ho}(\mathbf{r}) + g|\Psi|^2\right]\Psi = \mu\Psi$$

Blackboard:

Density distribution, chemical potential and cloud radius.

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Thomas	-Fermi appr	oximation		

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Thomas-Fermi approximation

Density distribution for $r \leq r_{\rm TF}$

$$n(\mathbf{r}) = \frac{\mu}{g} \left(1 - \frac{r^2}{r_{\rm TF}^2} \right)$$

Chemical potential

$$\mu = \frac{1}{2}\hbar\omega_{\rm ho} \left(15 \, \frac{Na}{a_{\rm ho}}\right)^{2/5}$$

Spatial extent of cloud

$$r_{\rm TF} = a_{\rm ho} \left(15 \, \frac{Na}{a_{\rm ho}}\right)^{1/5} > a_{\rm ho}$$



Figure : Density distribution in a harmonic trap.



Thomas-Fermi approximation



Figure : Density distribution of sodium atoms. From: Dalfovo et al. (1999)

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Healing	length			

Kinetic energy = interaction energy

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$$\frac{\hbar^2 k^2}{2m} = \frac{\hbar^2}{2m\xi^2} = g n(\mathbf{r})$$
$$\Rightarrow \quad \xi^2 = \frac{\hbar^2}{2mqn} = \frac{1}{8\pi an}$$

Solving the GPE for a box potential

$$\Psi(r)=\Psi_0 \tanh\left(\frac{r}{\sqrt{2}\xi}\right)$$

Length scale on which perturbations in the density distribution are "healed".



 $\label{eq:Figure:Density} \begin{array}{l} \mbox{Figure:Density distribution in a box potential} \\ \mbox{within Thomas-Fermi approximation (red) and} \\ \mbox{solution of the GPE (green)} \end{array}$



Figure : Vortices in a BEC. From: Australian Centre for Quantum-Atom Optics

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Conclusion	า			

- Ideal bose gas
 - Critical temperature
 - Condensate fraction
- Experimental techniques
 - Laser cooling
 - Evaporative cooling
- Weakly interacting bose gas
 - Parabola shaped density distribution
 - Spatial extent of the cloud
 - Healing length

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Thank you for your attention.