#### Martina Dombrowski (MartinaDombrowski@web.de)

In 1966 Mermin and Wagner proved the absence of long-range order in two-dimensional systems with a continuous symmetry (e.g. XY-model) for finite temperature T [1]. Though, Kosterlitz and Thouless showed the existence of a phase transition in 1972. A phase transition is produced by unbinding vortex-antivortex pairs at a critical temperature  $T_{KT}$  [2].

## 1 XY-Model

Two dimensional square lattice with planar rotors of unit length.



Figure 1: Schematic XY model

Find Hamiltonian H and interaction term J:

$$H = -J \sum_{\langle i,j \rangle} \mathbf{S}_i \cdot \mathbf{S}_j = -J \sum_{\langle i,j \rangle} \cos(\theta_i - \theta_j)$$
(1)

For low temperature with ground state energy  $E_0 = 2JN$ :

$$H = E_0 + \frac{J}{2} \int d^2 r (\nabla \theta)^2 \qquad (2)$$

Free energy F of a single vortex with system size L and lattice spacing a:

$$F = E_0 + (\pi J - 2k_B T) \ln(L/a)$$
 (3)

Indentify critical temperatur  $T_{KT} = \frac{\pi J}{2k_B}$ 

## 2 Phase Transition

#### 2.1 Mermin Wagner Theorem

Looking at the mean magnetisation  $\langle \mathbf{S} \rangle$ Mermin and Wagner showed that there is no long-range order to be found in 2D systems (e.g XY-model) for finite temperature.  $\langle \mathbf{S} \rangle$  goes to zero for finite temperature. Fluctuations at even low temperature will destroy the long-range order.

### 2.2 Correlation Function

Looking at the correlation function  $\langle \mathbf{S}(\mathbf{r})\mathbf{S}(\mathbf{0})\rangle$ , you will find quasi-long-range order:

– Low temperature:

$$\langle \mathbf{S}(\mathbf{r})\mathbf{S}(0)\rangle \simeq \left(\frac{r}{L}\right)^{-T/2\pi J}$$
 (4)

 $\rightarrow$  quasi-long-range order.

- High temperature

$$\langle \mathbf{S}(\mathbf{r})\mathbf{S}(0)\rangle = e^{-\frac{i}{\xi}}$$
 (5)

 $\rightarrow$  exponential decay.

Hauptseminar "Physik der kalten Gase", SS 2023, Univesität Stuttgart

### 2.3 Superfluid Density

Superfluid density  $\rho_s$  measures the energy change caused by a twist. Therefore it describes the effect of thermally activated vortex pairs.

$$\rho_s^R = \begin{cases} \rho_s^R (T_{KT}^-) [1 + const. (T_{KT} - T)^{1/2}] \\ 0 \end{cases}$$
(6)

with the first case showing  $T < T_{KT}$  and second case  $T > T_{KT}$ .



Figure 2: Superfluid density over T.

Universal for all systems you get  $\rho_s^R(T_{TK}^-)/T_{KT} = 2/\pi$ .

# 3 Kosterlitz-Thouless Phase Transition

Increasing the temperature activates vortices. For  $T < T_{KT}$  activated vortices bind to vortex-antivortex pairs. At  $T_{KT}$  all pairs unbind collective.

3.1 Experiment



Figure 3: Experiment with superfluid helium films.

## References

- [1] N.D. Mermin and H. Wagner, Phys. Rev. Lett. 22 1133 (1966).
- [2] J.M. Kosterlitz and D.J. Thouless, J. Phys. C 6, 1181 (1973).
- [3] P.M. Chaikin and T.C. Lubensky, *Principles of condensed matter physics*, Cambridge University Press (1995).
- [4] H. J. Jensen, *The Kosterlitz-Thouless Transition*, Department of Mathematics, Imperial College (2003).