

1.7. Topologically protected edge states in classical systems

1.7.1. Review: Effects of topological bands

* Quantized Hall response:

$$\sigma_{xy} = -\frac{e^2 v}{2\pi \hbar} \quad \text{with} \quad v = \sum_u C^{(u)} \quad \in \mathbb{Z}$$

↑
SP properties

$u: \epsilon_u < E_F$

MB quantum effect

→ Requires filled bands

→ Many-body phenomena (Fermi statistics)

→ Genuine quantum effect

* Robust edge states:

Top. system 1D (OTSC) \rightarrow zero-energy boundary modes

Top. system 1D (OTC) \rightarrow gapless edge modes

\rightarrow Bulk-boundary correspondence

\rightarrow Single-particle phenomena

\rightarrow Not a quantum effect \square

Topological features of the band structure (= single-particle feature)

are not quantum effects \square

Question: Can we translate edge modes to classical systems?

1.7.5. Example: Topological mechanics and helical edge modes

1. Goal: Realization of the helical edge modes of the QSHET in a classical mechanics setup governed by Newton's equations.

2. Quantum systems (QS): ∇ SP Schrödinger equations

$$\underbrace{i\hbar \dot{\psi}_i^{\alpha}}_{\text{Quantum dynamics}} = H_{ij}^{\alpha\beta} \psi_j^{\beta}$$

i, j : site indices,

α, β : spin indices

H : SP-Hamiltonian

Pauli coupled
by springs \hookrightarrow Hermitian matrix

3. Classical systems (CS):

∇ Newton's equation for N coupled oscillators

Remember:

$$m \ddot{x} = -Dx$$

$$\omega^2 = D/m$$

$$\underbrace{\ddot{x}_i}_{\text{Classical dynamics}} = -D_{ij} x_j$$

x_i : states of oscillators

D : Dynamical coupling matrix
(real, symmetric, positive semidefinite)

4. Observation:

- * QS characterized by eigenstates of H
- * CS characterized by eigenmodes of D

But: Edge modes are simply special eigenstates of H

Idea: Use H (with edge modes) to construct D (with edge modes).

5. Model construction:

a) Two independent copies of the Hofstadter model with spin-dependent flux $\alpha \hat{\phi} = \pm \frac{P}{q} = \pm \frac{1}{3}$

Landau gauge

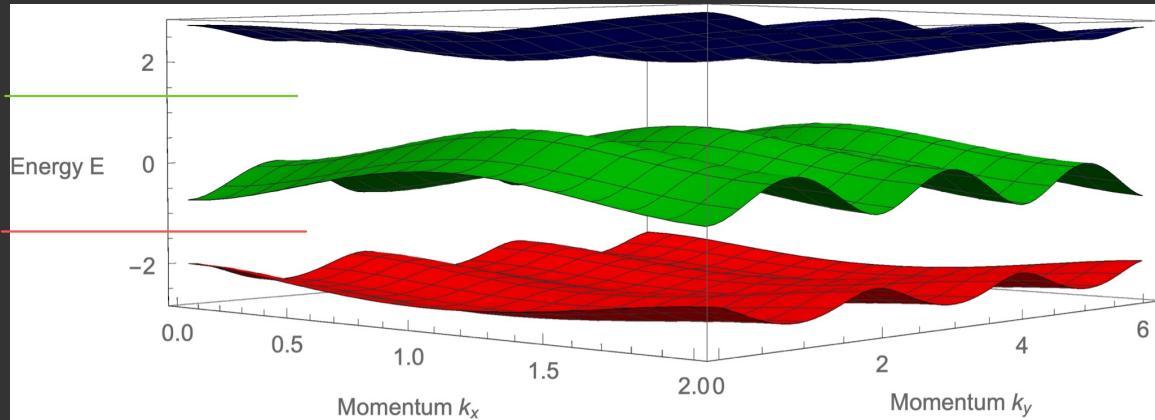
$$H = \begin{pmatrix} H_T & 0 \\ 0 & H_J \end{pmatrix} \quad \text{with} \quad H_\alpha = \sum_{x,y} \left\{ i[x, y+1, \alpha] X[x, y, \alpha] + e^{-2\pi i \alpha \hat{\phi} y} i[x+1, y, \alpha] X[x, y, \alpha] \right\} + h.c.$$

↳ $g=3 \rightarrow$ Three gapped, spin-degenerate bands

Chemical numbers

$$\Theta = C_1 + C_2 = -1$$

$$\Theta = C_1 = 1$$



$$\rightarrow C_3 = 1$$

$$\rightarrow C_2 = -2$$

$$\rightarrow C_1 = 1$$

↳ Finite sample with open boundaries:

→ Two helical edge modes in each gap

c) Symmetries:

* Time-reversal symmetry: $T = i\sigma^Y K$ with $T^2 = -1$

$$THT^{-1} = H$$

since $H_\uparrow = H^* \downarrow$

\rightarrow AII with \mathbb{Z}_2 Pfaffian index

* Spin conservation:

$$\sigma^z H \sigma^z = H$$

$$I = \frac{C_T - C_L}{2} \bmod 2$$

d) Problem: H is complex

$$H = H_\uparrow$$

Solution:

$$U = \underbrace{\frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -i \\ 1 & i \end{pmatrix}}_U \otimes \underbrace{\mathbb{1}_{\text{lattice}}}_{\mathbb{1}}$$

$$\Rightarrow D := U^\dagger H U = \begin{pmatrix} \operatorname{Re} H & \operatorname{Im} H \\ \operatorname{Im} H^\dagger & \operatorname{Re} H \end{pmatrix}$$

$$\operatorname{Re} H = \operatorname{Re} H_\uparrow = \operatorname{Re} H_L, \quad \operatorname{Im} H_\uparrow = \operatorname{Im} H_\uparrow^* = \operatorname{Im} H_\uparrow^\dagger = \operatorname{Im} H^\dagger$$

→ D real and symmetric (and can be made positive semidefinite by a constant shift $D \mapsto D + \text{const}$)

c) Transformed basis:

$$\begin{pmatrix} X_{xy} \\ Y_{xy} \end{pmatrix} = U^T \begin{pmatrix} \Psi_{xy\uparrow} \\ \Psi_{xyd} \end{pmatrix}$$

(X_{xy}, Y_{xy}) : Position of 2D harmonic oscillator on site (x, y)

Ψ spin-up mode on site (x, y) :

$$\Psi_{xy\uparrow}(+) = \Psi_0 e^{-i\omega t} \begin{pmatrix} 1 \\ 0 \end{pmatrix}_{xy} \leftarrow \text{"Spin Basis"}$$

$$\Rightarrow U^T \Psi_{xy\uparrow}(+) = \Psi_0 e^{-i\omega t} \underbrace{\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix}_{xy}}_{\text{Jones vectors}} \leftarrow \text{"Oscillator basis"}$$

$$= \text{Re}[U^T \Psi_{xy\uparrow}(+)] \propto \begin{pmatrix} \cos(\omega t) \\ \sin(\omega t) \end{pmatrix} \leftarrow \text{left-circular polarization}$$

Spin UP \rightarrow Left circular polarization

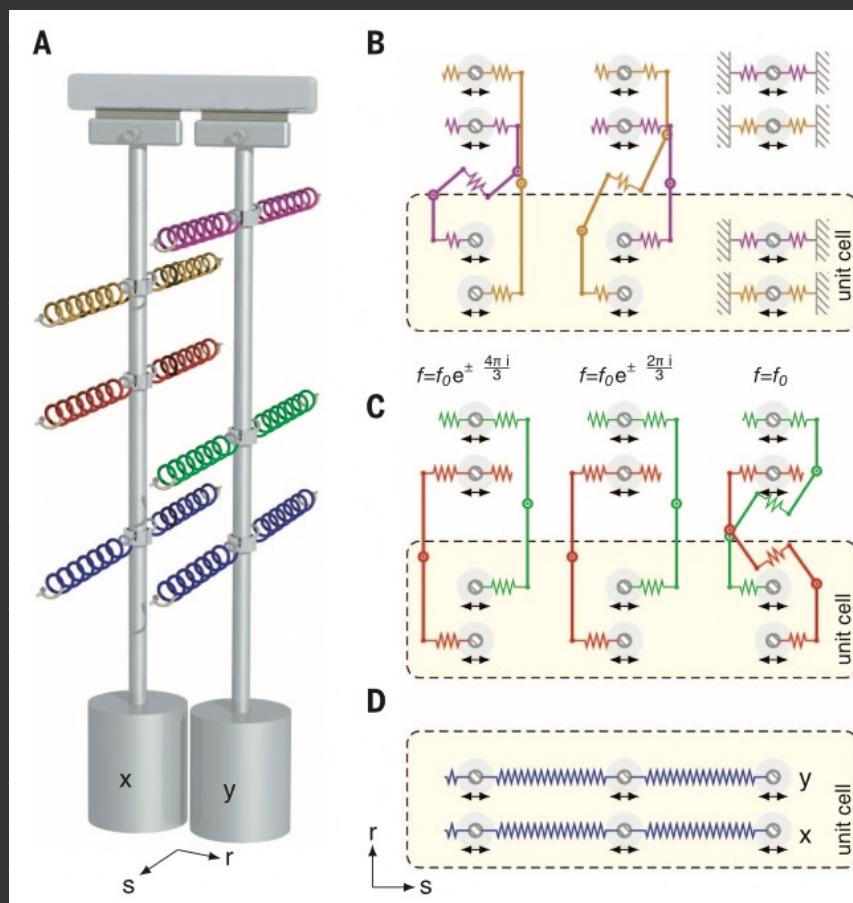
Spin DOWN \leftarrow Right circular polarization

f) Requirements on the setup:

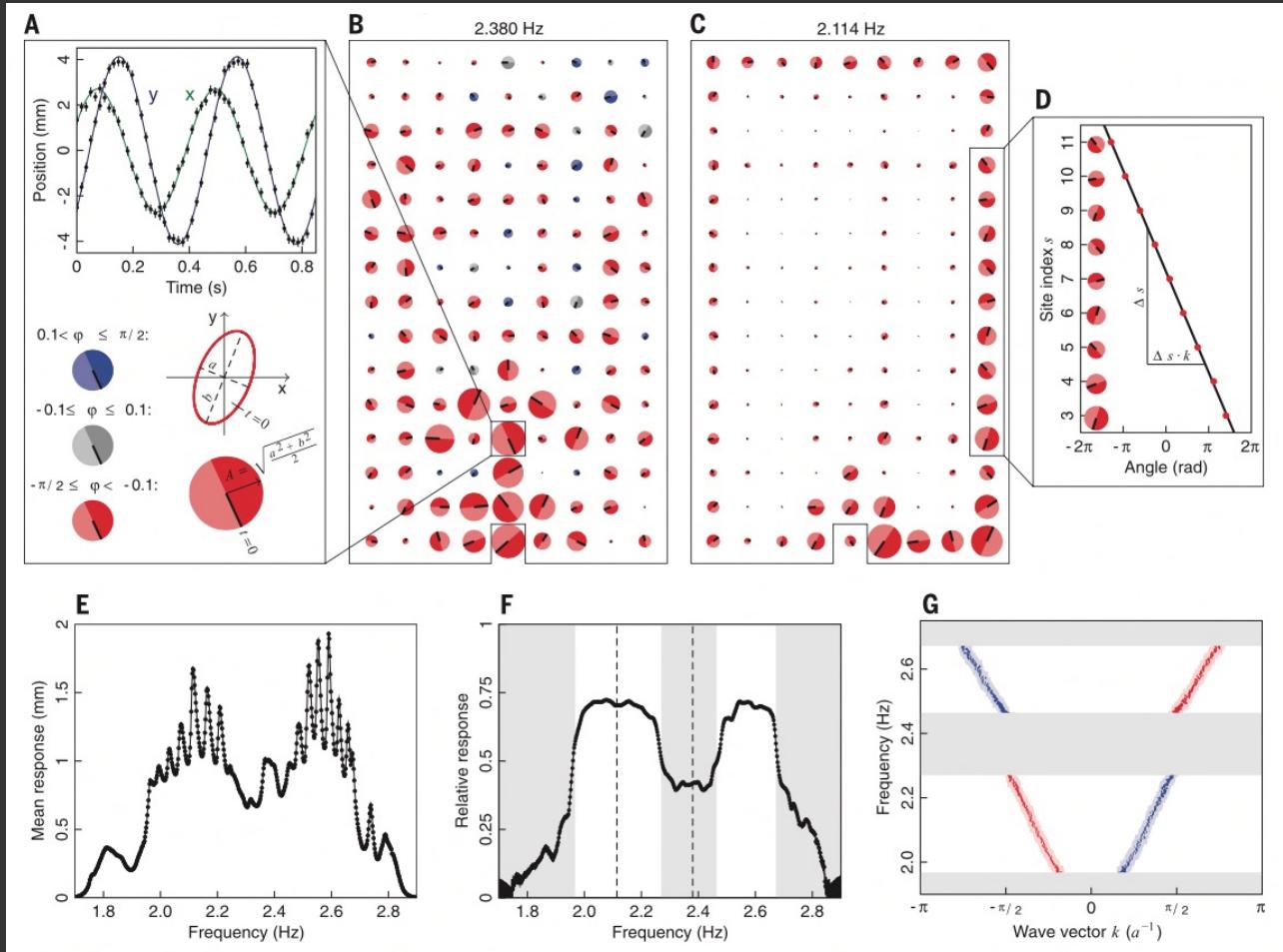
- * Nefstadsler model with $g=3 \rightarrow 3$ sites per magnetic unit cell
- * Two-fold spin/polarization degeneracy per site
 - \rightarrow 2 harmonic oscillators per site
 - $\rightarrow 6 \times 1D$ pendula per unit cell coupled by springs according to \mathcal{D} :
 - * $Re H \rightarrow XX$ - and YY -couplings
 - * $Im H \rightarrow XY$ -couplings

6. Experiment & Results:

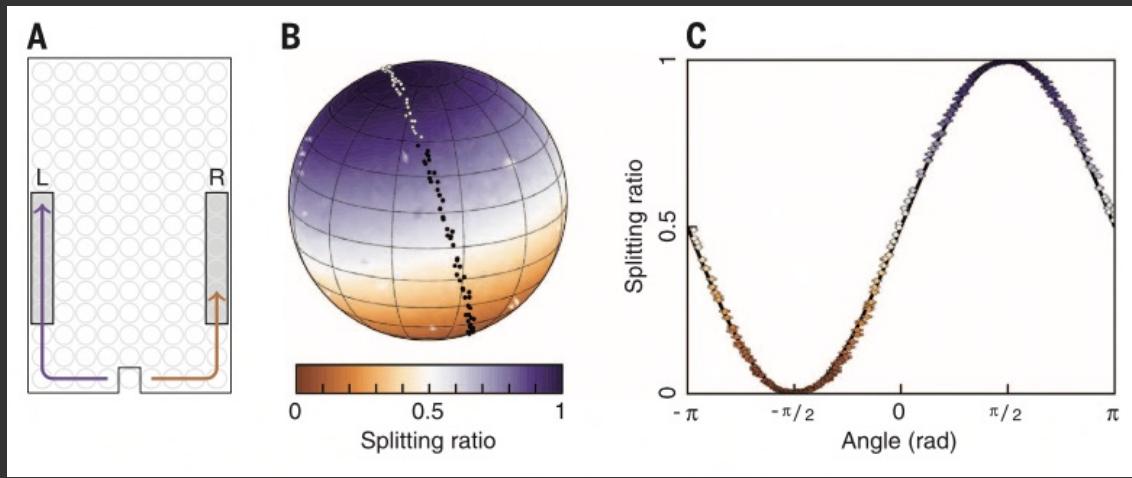
a) Construction:



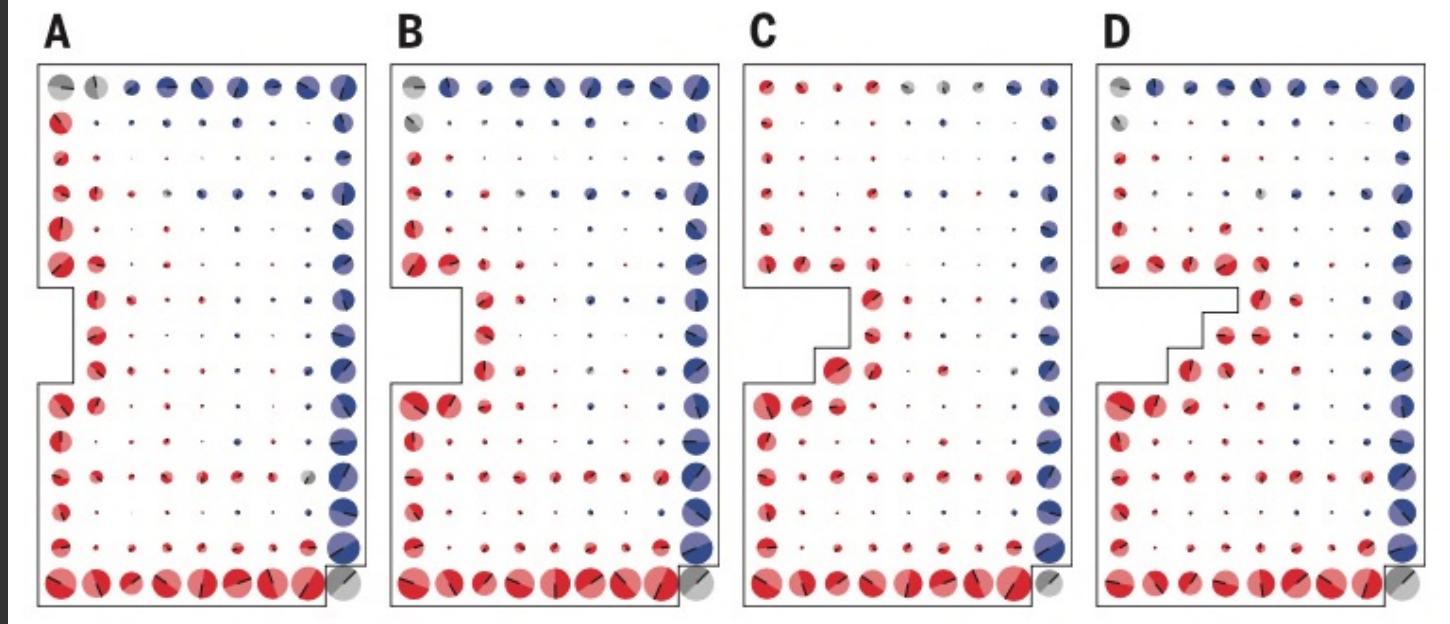
b) Edge modes:



c) Application: Beam splitter



d) Robustness



Symmetry protection:

i) Y

$$H \stackrel{!}{=} THT^{-1} = U_T H^* U_T^+ = U_T (UU^+ + UU^+)^* U_T^+$$

$$\Rightarrow D = U^+ H U = (U^+ U_T U^*) D \underbrace{(U^{*-1} U_T^+ U)}_S = S^+ D S$$

with unitary symmetry

$$S = U^{*-1} U_T^+ U = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

$$\rightarrow \text{with } D = \begin{pmatrix} D_{xx} & D_{xy} \\ D_{yx} & D_{yy} \end{pmatrix} \Rightarrow$$

$D_{xx} \stackrel{!}{=} D_{yy}$
 $D_{xy} \stackrel{!}{=} -D_{yx}$

Satisfied by $D_{xx} = D_{yy} = \text{Re } H$

$$D_{xy} = \text{Im } H$$

$$D_{yx} = \text{Im } H^T = \text{Im } H^* = -\text{Im } H = -D_{xy}$$

→ Postulation

$$D' = D + \begin{pmatrix} \delta D & 0 \\ 0 & \delta D \end{pmatrix}$$

does not destroy the edge modes!

→ Local symmetry constraint

1.7.3. More classical systems

- * Topological mechanics
- * Topological acoustics
- * Topological photons (→ Topological insulator lasers)
- * Topoelectrical circuits
- * Topological fluid dynamics

