

Problem 1: The Chern number as Skyrmion number
Learning objective

For the special case of 2-band Bloch Hamiltonians, the expression for the Chern number can be interpreted as a winding number that counts how often the sphere S^2 is wrapped when traversing the Brillouin zone. This number counts topological twists in the Bloch vector field on the Brillouin zone; these twists are known as (*Anti*-)Skyrmions. Here you derive the expression for the Chern number that is crucial for this interpretation.

Consider a translation invariant system in two dimensions with Bloch Hamiltonian

$$H(\mathbf{k}) = \vec{d}(\mathbf{k}) \cdot \vec{\sigma}. \quad (1)$$

Denote the ground and excited states as $|u_{\mathbf{k}}\rangle$ and $|v_{\mathbf{k}}\rangle$ with eigenenergies $-d(\mathbf{k}) = -|\vec{d}(\mathbf{k})|$ and $+d(\mathbf{k})$, respectively; the gap is therefore $2d(\mathbf{k})$.

The Chern number for Bloch bands was derived in the lecture and reads in this case

$$C = \frac{i}{2\pi} \int_{T^2} \left\{ \langle \tilde{\partial}_y u_{\mathbf{k}} | \tilde{\partial}_x u_{\mathbf{k}} \rangle - \langle \tilde{\partial}_x u_{\mathbf{k}} | \tilde{\partial}_y u_{\mathbf{k}} \rangle \right\} d^2k \quad (2)$$

with $\tilde{\partial}_i = \partial_{k_i}$.

Your goal is to massage this expression into a new form with a straightforward topological interpretation.

a) As a preliminary step, show that

$$\varepsilon_{ijk} \hat{d}_i (\tilde{\partial}_x \hat{d}_j) (\tilde{\partial}_y \hat{d}_k) = \varepsilon_{ijk} \frac{d_i (\tilde{\partial}_x d_j) (\tilde{\partial}_y d_k)}{d^3} \quad (3)$$

with $\hat{d} = \vec{d}/|\vec{d}|$.

b) Now show that

$$\langle v_{\mathbf{k}} | \tilde{\partial}_x u_{\mathbf{k}} \rangle = \frac{\langle v_{\mathbf{k}} | [\tilde{\partial}_x H] | u_{\mathbf{k}} \rangle}{-2d(\mathbf{k})} \quad (4)$$

and use this and the completeness of the Bloch basis $\{|v_{\mathbf{k}}\rangle, |u_{\mathbf{k}}\rangle\}$ to derive the expression

$$C = -\frac{1}{4\pi} \int_{T^2} \text{Im} \left[\langle u_{\mathbf{k}} | \sigma^i | v_{\mathbf{k}} \rangle \langle v_{\mathbf{k}} | \sigma^j | u_{\mathbf{k}} \rangle \right] \frac{\tilde{\partial}_y d_i \tilde{\partial}_x d_j}{d^2} d^2k. \quad (5)$$

(*Hint*: You may also use results from Problem 1 on Problem Set 2.)

c) Show that $\langle u_{\mathbf{k}} | \sigma^k | u_{\mathbf{k}} \rangle = -\hat{d}_k$ and with this

$$\text{Im} [\langle u_{\mathbf{k}} | \sigma^i | v_{\mathbf{k}} \rangle \langle v_{\mathbf{k}} | \sigma^j | u_{\mathbf{k}} \rangle] = -\varepsilon_{ijk} \hat{d}_k \tag{6}$$

so that with a) finally

$$C = -\frac{1}{4\pi} \int_{T^2} \hat{d} \cdot (\tilde{\partial}_x \hat{d} \times \tilde{\partial}_y \hat{d}) \, d^2k. \tag{7}$$

Identify the Berry curvature $\mathcal{F}_{xy}(\mathbf{k})$.

This is the expression used in the lecture to motivate the interpretation of Chern number and Berry curvature in terms of Skyrmion number and -density, respectively.

Problem 2: The Berry curvature of Dirac Hamiltonians

Learning objective

In the lecture we identified Dirac Hamiltonians as useful tools to study *changes* in Chern numbers. Here you derive a simple expression for the integral of the Berry curvature of a general Dirac Hamiltonian over its (non-compact) momentum space \mathbb{R}^2 .

Consider the general Dirac Hamiltonian

$$H_D(\mathbf{k}) = \sum_{i,j=1}^2 k_i h_{ij} \sigma^j + h_z \sigma^z \tag{8}$$

with 2×2 matrix h and “mass” h_z .

Show that the integral of the Berry curvature (of the lower band) over \mathbb{R}^2 yields

$$C = -\frac{1}{2\pi} \int_{\mathbb{R}^2} \mathcal{F}_{xy} \, d^2k = -\frac{\text{sign}[h_z] \text{sign}[\det(h)]}{2}. \tag{9}$$

(Hint: Use the result derived in Problem 1 and make the linear substitution $\mathbf{k}' = h^T \mathbf{k}$.)

Problem 3: Kramers theorem

Learning objective

Kramers theorem guarantees the degeneracy of eigenenergies for time-reversal invariant Hamiltonians of *half-integer total spin*. This has consequences in many situations, e.g., in atomic physics where it explains the degeneracy of energy levels for states with an odd number of fermions in the absence of a magnetic field. The theorem was mentioned in the lecture but not proven; here you complete this task.

Mathematically, Kramers theorem states that every eigenenergy of a time-reversal invariant ($[H, T_U] = 0$) Hamiltonian H with a (anti-unitary) time-reversal symmetry $T_U = UK$ that squares to minus one, $T_U^2 = -\mathbb{1}$, is at least two-fold degenerate.

Prove this!