

**Problem 1: Stark Effect for a Harmonic Oscillator (Oral)****Learning objective**

The goal of this problem is to apply a static perturbation (as derived in the lecture) up to second order for a simple setup. This setup has the special property that an exact solution exists and therefore the perturbation theory can be tested by a comparison with the exact solution.

The Hamiltonian of a one-dimensional harmonic oscillator within a homogeneous electric field  $E$  is given by

$$H = \frac{1}{2} (P^2 + Q^2) + eEQ, \quad (1)$$

where the Hamiltonian is dimensionless, i.e.,  $[Q, P] = i$ . Consider the second term of the Hamiltonian as a perturbation of the free oscillator, that is,  $H_1 = Q$  and  $\lambda = eE$ .

- Calculate the perturbed eigenfunctions and energy eigenvalues up to second order in  $\lambda$ .
- Compare the result of perturbation theory with the exact solution for the energy of the problem.

**Problem 2: Fermi's Golden Rule (Written)****Learning objective**

This assignment is a repetition of the derivation of Fermi's golden rule. To this end, a system with a periodic perturbation in time that is adiabatically switched on is studied; this requires some modifications of the proof presented in the lecture.

Consider a system with Hamiltonian  $H_0$ , eigenstates  $|n\rangle$  and eigenenergies  $E_n$ . We prepare the system in the initial state  $|n = i\rangle$  at  $t_0 = -\infty$  and slowly ramp up the periodic perturbation  $V \cos(\omega t)$  with frequency  $\omega$  and static perturbation  $V$ , i.e.,

$$H'(t) = V \cos(\omega t) \exp(\eta t) \quad (2)$$

for  $\eta > 0$  (so that  $H'(t) = 0$  for  $t \rightarrow -\infty = t_0$ ).

Derive Fermi's golden rule in the adiabatic limit  $\eta \rightarrow 0$ , i.e., calculate the transition rate  $\Gamma_{i \rightarrow f}$  from the initial state  $|i\rangle$  to some final state  $|f\rangle$  with energy  $E_f$ .

**Hint:** Use the Dirac sequence  $\lim_{\eta \rightarrow 0} \frac{\eta}{x^2 + \eta^2} = \pi \delta(x)$  and neglect the oscillating contributions to the transition rate (why is this reasonable?).

**Problem 3: Time-Dependent Perturbation Theory (Written)**

**Learning objective**

The first subtask is an application of time-dependent perturbation theory in first order. The goal is to identify the correct transition amplitude and evaluate the integrations analytically. The second subtask is a challenging calculation to derive the *exact* solution of the problem—a useful repetition/application of important methods in quantum mechanics. This problem is one of the few examples where the time evolution can be solved analytically (which allows for a verification of perturbation theory).

We investigate a one-dimensional harmonic oscillator with mass  $m$ , charge  $e$ , and frequency  $\omega$  in a time-dependent electric field  $E(t)$ . The Hamiltonian is of the form

$$\begin{aligned}
 H &= H_0 + H'(t), \\
 \text{where } H_0 &= \frac{p^2}{2m} + \frac{m}{2}\omega^2 x^2 \quad (\text{harmonic oscillator}) \\
 \text{and } H'(t) &= ex E(t) \quad (\text{perturbation}).
 \end{aligned} \tag{3}$$

The time-dependency of the external electric field is given by

$$E(t) = \frac{A}{\tau\sqrt{\pi}} e^{-(t/\tau)^2}, \tag{4}$$

where  $A \in \mathbb{R}$  is a constant and  $\tau > 0$  is the decay rate.

- a) Calculate the transition probability  $P_{0 \rightarrow n}(t, t_0)$  from the ground state  $|0\rangle$  at  $t_0 = -\infty$  to an excited state  $|n\rangle$  at  $t = +\infty$  in first order perturbation theory. What happens for  $\tau \rightarrow 0$ ?

**Hint:** Use  $x = \sqrt{\frac{\hbar}{2m\omega}} (a^\dagger + a)$  to evaluate the matrix element.

- b) The transition probability can also be calculated *exactly*.

Prove that the exact expression of  $P_{0 \rightarrow n}(t, t_0)$  for  $t_0 = -\infty$  and  $t = +\infty$  reads

$$P_{0 \rightarrow n} = \frac{K^{2n}}{n!} e^{-K^2} \quad \text{with} \quad K = \frac{eA}{\sqrt{2m\omega\hbar}} e^{-\omega^2\tau^2/4}, \tag{5}$$

and compare the result with a).

**Hint:** Use the Dirac picture (as introduced in the lecture) and derive the equation of motion

$$i\hbar\partial_t |\psi_D(t)\rangle = H'_D(t) |\psi_D(t)\rangle \quad \text{with} \quad H'_D(t) = \sqrt{\frac{\hbar}{2m\omega}} eE(t) [e^{-i\omega t} a + e^{i\omega t} a^\dagger]. \tag{6}$$

Then use the ansatz  $|\psi_D(t)\rangle = \exp[-iK(t)a^\dagger] |\bar{\psi}(t)\rangle$  and choose  $K(t)$  as

$$K(t) = \frac{e}{\sqrt{2m\hbar\omega}} \int_{-\infty}^t e^{i\omega s} E(s) ds \tag{7}$$

to eliminate  $a^\dagger$  from the Schrödinger equation. To this end, the relation  $e^{iK(t)a^\dagger} a e^{-iK(t)a^\dagger} = a - iK(t)$  can be used to derive a differential equation for  $|\bar{\psi}(t)\rangle$ . Solve this equation with an ansatz  $|\bar{\psi}(t)\rangle \propto |0\rangle$  to find the exact solution  $|\psi(t)\rangle$  for the wave function in the Schrödinger picture. The transition probability is then  $P_{0 \rightarrow n} = |\langle n|\psi(t = \infty)\rangle|^2$ .